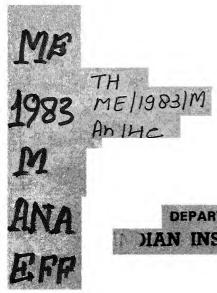
EFFECT OF TITATANIUM, ZIRCONIUM AND TIN ON THE VARIATION OF SATURATION MAGNETISATION, CURIE TEMPERATURE AND LATTICE PARAMETER OF Ni.3 Zn.7 Fe<sub>2</sub> O<sub>4</sub>

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DEPARTMENT OF METALLURGICAL ENGINEERING

DIAN INSTITUTE OF TECHNOLOGY KANPUR

AUGUST, 1983

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V. S. ANANTHAN

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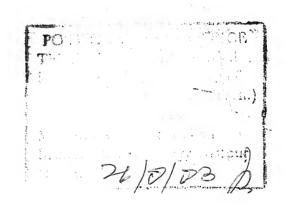
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## CERTIFICATE

Certified that the thesis entitled, 'EFFECT OF TITANIUM, ZIRCONIUM AND TIN ON THE VARIATION OF SATURATION MAGNETISATION, CURIE TEMPERATURE AND LATTICE PARAMETER OF Ni 3Zn 7Fe2O4 FERRITE: has been carried out under our supervision and it has not been submitted elsewhere for a degree.

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#### CHAPTER 1

#### INTRODUCTION

The outstanding feature of ferrites as magnetic materials is their high resistivity which makes them particularly suitable for applications in high frequency systems where its high resistivity ensures low losses. The ferrite industry was established about thirty-five years ago following the classical investigations of J.L. Snoek<sup>1</sup>.

Ferrites are oxides with a formula MeOFe<sub>2</sub>O<sub>3</sub>, where Me represents divalent cations like Mn,Fe,Co,Ni,Cu,Zn etc., and has aninverse spinel structure. Oxygen with relatively larger radii form an face centered cubic lattice. In this closed packed structure two kinds of interstitial sites occur, the tetrahedral (A) and octahedral (B), sites which are surrounded by 4 and 6 oxygen ions respectively. A unit cell has 32 O<sup>2-</sup> ions with 64 tetrahedral and 32 octahedral sites. The metal ions having radii 0.4 - 1.0 A are distributed amongst these sites.

The distribution of cations at A and B interstices give rise to three type of spinel structures:

(1) 'Normal spinel' in which Me occupies the tetrahedral position and Fe<sup>3+</sup> in the octahedral position and is written as Me [Fe<sub>2</sub>]O<sub>4</sub>.

- (2) 'Inverse spinel' in which Me occupies the octahedral position together with half of Fe<sup>3+</sup> while the other half Fe<sup>3+</sup> occupies the tetrahedral position and is written as Fe[MeFe]O<sub>A</sub>.
- (3) 'Intermediate spinel' as Fe<sub>1-x</sub>Me<sub>x</sub>[Me<sub>1-x</sub>Fe<sub>1+x</sub>]O<sub>4</sub>.

The distribution of cations amongst these sites based on crystal field theory was explained by Gorter<sup>2</sup>. According to this theory, the magnetic interaction (super-exchange) energy is strongly dependant on the distribution of magnetic ions amongst the crystallographic positions. The influence of diameter and charge of a cation on its preference to A or B site in an oxidic spinel is not an individual property as the diameter and charge of other cations must be taken into account.

Later, Blasse<sup>3</sup> deduced the site preference energy qualitatively using Ligand field theory. Ligand field theory is the use of molecular orbitals and introduces partly covalent bonding and crystal field stabilization simultaneously. The results obtained are more reliable than those due to crystal field theory.

Ni<sup>2+</sup> has a strong preference for 6-fold co-ordination and Zn<sup>2+</sup> for 4-fold co-ordination. The preference of zinc to A site is utilised to improve the property of nickel ferrite. The addition of Zn<sup>2+</sup> ions to nickel ferrite forces

Fe<sup>3+</sup> ions from A to B sites. This should give an addition magnetisation of B-sites. However, Gorter<sup>2</sup> found that it does not increase. The addition of more than 50 mole perce of zinc in nickel ferrite reduces the net magnetic moment from the peak value as the reduced number of Fe<sup>3+</sup> ions in A sites becomes less able to maintain the alignment of B sublattice moments against BB interaction. The curie temperature  $(T_c)$  decreases linearly with increase of zinc content.

The variation of saturation moment, curie temperature and susceptability of ferrimagnetic compounds as a result of substitution of nonmagnetic ions in place of magnetic ions can be understood to a good approximation on the basis of random distribution of incomplete superexchange interactions. Gilleo<sup>4</sup> has derived expressions for saturation magnetisation and curie temperature based on superexchange interaction of magnetic ions of different co-ordination. Using Gilleo's formula, curie temperature for Cu<sub>x</sub> Zn<sub>1-x</sub> Fe<sub>2</sub>O<sub>4</sub> samples were evaluated by Swant et al.<sup>5</sup>

The exchange interactions between magnetic ions in A and B sites govern the magnetic ordering temperature  $(I_c)$ . The interaction energy between a magnetic moment 'S<sub>i</sub>' on site 'i' and 'S<sub>i</sub>' on site 'j' is given by the expression

Eexchange = -2Jij Si Sj

where  $J_{ij}$  is the exchange constant which depends on the distance between magnetic ions and the symmetry of local arrangement of anions and cations. The various interactions possible in a spinel are AA,BB and AB interactions.

Bongers et al.<sup>6</sup> studied the curie temperature variation as a function of  $Fe^{3+}$  on A and B sites for Rh, Ti and Sb substituted MgFe<sub>2</sub>O<sub>4</sub> and calculated  $J_{AB}(Fe^{3+}-Fe^{3+})$  using the slope of the curve. He was also able to account this temperature variation on the basis of strong BB interaction

The exchange constants for titanium substituted Ni 3 Zn .7 Fe<sub>2</sub>O<sub>4</sub> system using 3-sublattice Yafet-Kittel type model of Satayamurthy et al. 7, was calculated by Misra<sup>8</sup> and found the results agreeable with those of Sreevastava et al. 9.

Gorter<sup>2</sup> reported that the effect of Ti<sup>4+</sup> ion in an inverse spinel ferrite is to replace the ferric ion completely or partly thereby reducing the saturation moment. He had worked on the effect of Ti<sup>4+</sup> ion on nickel ferrite and nickel zinc ferrite where the titanium ion occupies the octahedral sites only.

According to J.E. Knowles<sup>10</sup> the effect of Ti<sup>4+</sup> ion in Mn-Zn ferrite is to localize ferric ion which increases resistivity and reduces loss. The anomalous results when substituted by Sn<sup>4+</sup> ion was attributed to the large size of

tin ion. When ferrite is cooled down from high temperature the quatrivalent ions tend to migrate to the preferred B sites which is facilitated by the presence of cation vacancies. During this process up to 1/3 of larger tin ion gets 'frozen in' in A site thereby increasing the magnetisation. But the replacement of  $Ti^{4+}$  ion by  $Sn^{4+}$  ion lowers curie temperature  $(T_c)$ . The lattice parameter is also found to increase with tin content.

Stinges et al.  $^{11}$  found the substitution of  $T_i^{4+}$  in Mn-Zn ferrite takes place on octahedral sites by replacing  $2Fe^{3+}$  by  $Fe^{2+}+Ti^{4+}$  under reducing conditions, which increases magnetic anisotropy. The lowering of curie temperature with titanium addition was found to be very small. He had also studied the variation of saturation magnetisation ( $M_s$ ) with titanium addition.

For Ni<sub>1-x</sub> Zn<sub>x</sub> Fe<sub>2</sub>O<sub>4</sub> as x value decreases the material becomes magnetically softer to harder. In soft ferrite as in our Ni <sub>3</sub> Zn <sub>7</sub> Fe<sub>2</sub>O<sub>4</sub>, the internal magnetic forces are relatively small, implying that anisotropy is small and domain walls movement are easier. Thus the magnetic spins yield easily to external field giving high permeability. The coercivity is small and the curie point is low.

The properties of ferrites depend on microstructure and chemical composition. The effect of processing parameters on the microstructure development of Ni. 3 Zn. 7 Fe<sub>2</sub>O<sub>4</sub> ferrite is discussed by Gupta<sup>12</sup>. The properties are sensitive to stoichiometry<sup>13</sup>. Sintering of such ferrites are carried out at 1100-1250°C. High sintering temperature leads to dense product. Slick and Blassches<sup>14</sup> noted sublimation of zinc oxide in Mn-Zn ferrite as a result of decrease of oxygen content or increase of sintering temperature. The partial pressure of zinc is maintained by the use of packing material while firing thereby avoiding the loss of zinc from the system.

V.V. Pankov et al. 15 have worked on the mechanism of Ni-Zn ferrite formation. They have shown that the reaction zone of interacting NiO, ZnO with Fe<sub>2</sub>O<sub>3</sub>, the ferrite phase crystallizes on iron oxide. Also, the interaction of (Ni, Zn)O solid solution with Fe<sub>2</sub>O<sub>3</sub> takes place by the mechanism Fe<sup>3+</sup>, Fe<sup>2+</sup> and Ni<sup>2+</sup>, Zn<sup>2+</sup> ion interdiffusion and have applied Wagner-model to determine the reaction rate constant.

Miter T. Dimova<sup>16</sup> found that in Ni-Zn ferrite the concentration of free Fe<sub>2</sub>O<sub>3</sub> was 1/3 to 1/2 after double ferritisation compared to the single process under same conditions.

The lattice constant of the whole range of  $Ni_x$   $In_{1-x}$   $Fe_2O_4$  were studied by Guilland IT and the value is adopted in

A.S.T.M<sup>18</sup>. The variation of saturation magnetisation with temperature for some ferrites of  $Ni_{1-\delta}$   $Zn_{\delta}$   $Fe_2O_4$  series are given in Smit and Wijn<sup>19</sup>. The saturation magnetisation, an important property of soft ferrites, was extensively studied by  $Gorter^2$  and brought out a clear picture in the magnetic saturation moment between theoretical and practical values.

#### CHAPTER 2

#### AIM OF THE PROBLEM

High-valent non-magnetic cations when substituted in ferrite are expected to create vacancies under oxidising conditions. This should result in a monotonic variation of lattice parameter, saturation magnetisation and curie temperature. However, a sharp initial dip and subsequent rise in the lattice parameter, curie temperature and saturation magnetisation was observed in case of titanium substituted Ni-Zn ferrite by earlier workers <sup>8,25</sup>. Similar behaviour of magnetic moment in other high-valent substitutions was also observed <sup>25</sup>. The samples prepared by the above workers were by solid-state reaction and the chemical homogenity of the final material were thought to be insufficient.

Hence, it was decided to add high-valent dopants such as titanium, zirconium and tin by solution method to Ni $_3$ Zn $_7$ Fe $_2$ O $_4$  and verify the following parameters.

- 1) Variation of lattice parameters,
- 2) Variation of saturation magnetic moment, and

3) Curie temperature variation with the addition of the dopants.

The dopants  ${\rm Ti}^{4+}$ ,  ${\rm Zr}^{4+}$  and  ${\rm Sn}^{4+}$  have ionic radii close to that of  ${\rm Ni}^{2+}$ ,  ${\rm Zn}^{2+}$  and  ${\rm Fe}^{3+}$  so that they can substitute ions from base Ni-Zn ferrite.

## CHAPTER 3

## EXPERIMENTAL METHOD

## 3.1 RAW MATERIALS

The raw materials used for preparation of Ni-Zn-ferrite and corresponding dopants are given below.

Raw materials	Grade
a) Nickel metal (Ni)	> 99%
b) Zinc metal (Zn)	> 99%
c) Iron metal (Fe)	Crystal bar grade > 99.5%
d) Titanium metal (Ti)	Crystal bar grade > 99.5%
e) Zirconium metal (Zr)	Crystal bar grade >99.5%
f) Tin metal (Sn)	>99%
g) Sulphuric acid (H <sub>2</sub> SO <sub>4</sub> )	$A_{\bullet}R_{\bullet}$
h) Nitric acid $(H^{NO}_3)$	A.R.
i) Hydrofluoric acid (HF)	L.R.

For packing material iron oxide, nickel oxide and zinc oxide of L.R. grade are used.

## 3.2 PREPARATION OF THE PACKING MATERIAL

The nickel oxide, zinc oxide and iron oxide powders were weighed accurately in an electronic pan balance

(Osbar, Germany) following Table 2. Hand pellets were made with 5 percent PVA (polyvinyl alcohol) as the binder. The pellets were fired directly in an electrically heated silicon-carbide furnace at 1200°C for 6 hours in sillimanite crucibles. After firing, the pellets were ground in an agate mortar to fine powder (-50 mesh).

## 3.3 PREPARATION OF SAMPLES

## 3.3.1 Preparation of base solution

Nickel: The nickel metal powder was weighed accurately as in Table 2. The metal was dissolved in hot concentrated nitric acid. The required amount of the solution was made up.

Zinc: The zinc metal was cut into small pieces with diamond wheel cutter and weighed accurately. It was then reacted with hot concentrated sulphuric acid. The zinc sulphate formed was slowly dissolved in distilled water by constant stirring and slow heating. The required amount of solution was prepared.

Iron: The weighed amount of iron was dissolved in 50 percent concentrated nitric acid and the required amount of solution was prepared.

## 3.3.2 Mixing of the base solution

The base solutions of nickel, zinc and iron containing their respective salts were then thoroughly mixed to give a clear solutions. The amount of solution prepared was noted and the strength was calculated.

## 3.4 PREPARATION OF DOPANT SOLUTIONS

Titanium: Known amount of titanium (Table 3) was taken in a platinum crucible and hydrofluoric acid was added. It was gently heated to completely dissolve the metal. The solution was then added to a beaker containing sulphuric acid and a predetermined amount of solution was prepared.

Zirconium: In a similar manner as titanium, zirconium was treated with hydrofluoric acid for 6 to 8 hours to completely dissolve the metal. This solution was then added to sulphuric acid to prepare the standard solution.

Tin: Tin was heated in concentrated sulphuric acid when the metal was completely dissolved and a predetermined amount of solution was prepared.

## 3.5 ADDITION OF DOPANT SOLUTIONS

The exact volume of dopant solutions required for adding to the base solution in order to get the required weight percentages of the dopants were calculated. The calculated volumes of the dopant solutions were added to the base solution slowly with uniform stirring.

Dopant-added solutions of different percentages were separately prepared as shown in Table 3.

## 3.6 DRYING OF THE MIXTURE OF SOLUTIONS

The solutions prepared as above were then heated to dryness in an oven kept in the fuming chamber. The heating time was about 4 to 5 hours.

## 3.7 THERMAL DECOMPOSITION

The dried salts were than transferred to alumina crucibles. Then they were heated in an electric furnace at  $800^{\circ}$ C for 4 hours. At this temperature the decomposition of all the salts present takes place to form the corresponding oxides. A fine mixture of oxides was got in correct propositions of weight.

## 3.8 HEAT TREATMENT

The decomposed sample in the form of fine powder was ground in an agate mortar to yield a thorough mixture. With the use of small amount of PVA as the binder, pellets were hand formed. These pellets were covered with packing material and fired at 1200°C for 2 hours.

After cooling, the samples were reground to fine powder. Small pellets of 3 mm diameter were made in a small punch and dye with the help of PVA as binder. The remaining sample was again hand formed with the same binder. The pellets were refired at 1200°C for 4 hours in alumina crucibles.

After cooling, the hand formed samples were ground to fine powder (-50 mesh) for x-ray studies. The 3 mm diamter pellets were used for magnetic studies.

Throughout the firing operation, the temperature of the furnace was controlled carefully by the use of on-off controller (Leeds and Northrup, Electromax). A temperature profile of the heating muffle was drawn with the help of Platinium-Platinum 10 percent Rhodium thermocouple to locate the region of desired temperature.

#### 3.9 X-RAY ANALYSIS

## 3.9.1 Optimising the ferritisation temperature

The base Ni 3 Zn 7 Fe<sub>2</sub>O<sub>4</sub> sample powder was studied in Seifert X-ray diffractometer (Germany) with chromium as target. X-ray diffractograms were taken for samples fired at different temperatures and the data were compared with ASTM data of the NiO, ZnO and Fe<sub>2</sub>O<sub>3</sub> and for the final product.

## 3.9.2 Characterisation of Ni-Zn ferrite

The X-ray diffractogram of the base sample was taken and the d-spacing was compared with ASTM data for Ni-Zn ferrite.

## 3.9.3 Step scan analysis

The shift in d-values for different dopants at various weight percentages can be found out only by noting

the exact peak positions. This was done by the use of automatic step scan available in Seifert X-ray diffractometer.

First by knowing the approximate peak positions, the sample was set about 0.5 degree less than the maximum peak position. The angles were then changed by steps of 0.05 degrees and the number of counts for ten seconds at each step were printed automatically. The counts were plotted against the angle to determine the exact peak position. This was repeated for all peaks and samples.

## 3.9.4 Lattice Parameter Calculations

The measured angle (20) values and the wavelength of chromium target used were given as data in the computer program for calculating the d value and the corresponding lattice parameter. The lattice parameter values were plotted against the Nelson-Riley function  $f(\theta)$ . The curve of  $f(\theta)$  versus the lattice parameter (a) is expected to be a straight line and the extrapolated 'a' corresponding to  $f(\theta) = 0$ , using the least square fit gives accurate value of 'a' (Refer Appendix).

#### 3.10 MAGNETIC MEASUREMENTS

## 3.10.1 Saturation magnetisation

The samples were prepared for magnetic studies as in Section 3.8. The pellets were weighed accurately. The

samples were then tested in the magnetic field of a parallel field vibrating magnetometer (PARC model 150A). In this magnetometer, the samples are kept at vibration in vertical direction and the magnetic field applied acts horizontally. The magnet is capable of producing a maximum field of 10 K Oe with a pole gap of 2.5 inches. The induced a.c field produced by the sample in a pair of secondary coils placed on both sides of the sample is amplified and compared with the signal produced by a standard magnet giving rise to an output signal which is proportional to the magnetic moment of the sample.

The magnetic field was applied gradually in steps upto 10 KOe, and the saturation moment of all samples at this field were taken at room temperature.

The saturation moment is given by

$$\sigma = \sigma_0/m \text{ (emu/gm)} \text{ or by } 4\pi M_s \text{(gauss)}$$

where  $\sigma_0$  is the observed magnetic moment, m is the mass of the sample, and M<sub>S</sub> is given by M<sub>S</sub> =  $\sigma$  x , ( -density gms/cc). 3.10.2 Curie temperature

The variation of magnetic moment with temperature at residual field of 40 Oe were observed for all the samples. The temperature of the samples were measured by NiCr-NiAl thermocouple. The heating rate of the samples were slow and

uniform. Near the expected value of curie temperature ( $T_c$ ). the magnetic moment variation at very small (2-5 degree centigrade) were noted. Then  $\frac{d\sigma}{dT}$  versus T curves were plotted using interpolation techniques. The ferromagnetic transition temperature ( $T_c$ ) is defined as that temperature at which  $\frac{d\sigma}{dT}$  versus T curve shows a minimum.

#### CHAPTER 4

#### RESULTS

#### 4.1 RAW MATERIALS

The samples and the packing material were prepared in accordance with the compositions given in Tables 2 and 3. About 10-15 grams of batch of sample and 500 grams of packing material were prepared.

#### 4.2 PREPARATION OF PACKING MATERIAL

A ferritised powder was obtained when the raw materials were fired at  $1200^{\circ}\text{C}$  for 6 hours. The composition was chosen similar to the basic Ni  $_3\text{Zn}$   $_7\text{Fe}_2\text{O}_4$  ferrite.

Zinc has a low vapour pressure and sublimates at high temperatures. To prevent the expulsion of zinc from the sample during firing, proper oxygen partial pressure was provided by covering the samples with packing material. Pure zinc-oxide was not used as the packing material to avoid the diffusion of zinc from the highly concentrated packing zone into the sample.

## 4.3 PREPARATION OF SAMPLES

## 4.3.1 Preparation of base-solutions

Nickel: The nickel metal was readily dissolved in hot concentrated nitric acid to form nickel nitrate. The

nickel nitrate formed decomposes to nickel oxide before it becomes anhydrous 20 as

$$2Ni(NO_3)_2 = 2NiO + 4NO_2 + O_2$$

Zinc: Zinc metal forms zinc sulphate on heating with concentrated sulphuric acid. Zinc sulphate decomposes at 600°C to form zinc oxide<sup>21</sup>.

Iron: The iron metal was dissolved in nitric acid with slight heating. Iron nitrate decomposes to iron oxide on heating  $^{22}$  as

Fe + 
$$2HNO_3$$
 =  $Fe_2(NO_3)_2 + H_2$   
 $6Fe(NO_3)_2 + 5H_2O = 3Fe_2O_3 + (2NO+10HNO_3)$ 

## 4.3.2 Mixing of the base-solutions

The mixture of solutions containing the salts of iron, nickel and zinc was clear and homogenous.

## 4.4 Preparation of dopant solutions

The dopants were dissolved in acid and then added to the base-solution. This was done to improve the homogenity of the dopants with the basic Ni-Zn ferrite. Titanium and zircanium did not dissolve in sulphuric acid as the metals were not in the powdered form. So these metals were initially dissolved in hydrofluoric acid and further treated with sulphuric acid to form the corresponding metal sulphates 23,24. The solutions were gently heated to evaporate the hydrofluoric acid. The evaporation of hydrofluoric acid was observed as white fumes. The titanium and zirconium sulphates decompose at 708°C and 798°C respectively 25. Tin dissolved readily in sulphuric acid to form tin sulphate which decomposes to tin oxide on heating.

## 4.5 ADDITION OF DOPANT SOLUTIONS

The amount of dopant solutions added to the base-solution was in accordance to Table 3. The solutions were added using a micro-burette which reads up to 0.05 ml.

#### 4.6 DRYING

The solutions when dried gave a mixture of salts containing iron, nickel, zinc and dopants in proper propositions.

#### 4.7 THERMAL DECOMPOSITION

The salts of various elements present after drying were thermally decomposed to form their corresponding oxides. The complete decomposition was verified by their weight loss. The oxides were obtained in the form of fine particles.

## 4.8 HEAT TREATMENT

The initial grinding of the metal oxides improved the homogenity. The handformed pellets were made to improve ferritisation rate. The use of packing material avoided zinc escaping from the system.

After initial firing at 1200°C, the material was reground thereby further improving the homogenity of the metal oxides. Completely ferritised sample was obtained after refiring at 1200°C for 4 hours.

## 4.9 X-RAY ANALYSIS

## 4.9.1 Optimising the ferritisation temperature

The X-ray diffractograms were taken for all Ni<sub>3</sub>Zn<sub>.7</sub>Fe<sub>2</sub>O<sub>4</sub> samples fired at different temperatures, ranging from 1000 to 1150°C. The ferritisation was complete at 1150°C. This was indicated by the peaks in the diffractogram when they were identified with that of the final product. The firing temperature was chosen 50°C above the required temperature to ensure complete ferritisation.

## 4.9.2 Characterisation of Ni-Zn ferrite

The result of the X-ray diffraction lines of base Ni<sub>.3</sub>Zn<sub>.7</sub>Fe<sub>2</sub>O<sub>4</sub> is given in Table 4. The d-spacings were compared with the ASTM data for Ni-Zn ferrite. The hkl

indices were adopted from the ASTM data. The (311) planes refer to the maximum observed intensity.

## 4.9.3 Step scan analysis

The counts of the peaks in steps of 0.05 degree were taken. The angle (20) versus the counts were plotted and the curves in all the cases were smooth. Hence, the exact 20 value was determined as given in Appendix.

## 4.9.4 Lattice parameter

Accurate determination of lattice parameter is discussed in Appendix. The results of estimated d, hkl, lattice parameter and the corresponding Nelson-Riley functions are indicated in Tables 5,6 and 7. The extrapolated values of lattice parameter of all the samples are given in Tables 5,6,7 and 9.

Figure 1 shows the variation of lattice parameter with the addition of titanium, zirconium and tin in Ni  $3^{Zn} \cdot 7^{Fe} 2^{O} 4^{\bullet}$ 

It is observed for very small additions of titanium the lattice parameter decreases and for 0.5 wt% and above the lattice parameter increases. The dip is found minimum at 0.25 wt% of titanium when the lattice parameter decreases from 8.4178A to 8.4127A. This is followed by a continuous increase at higher percentages as seen in Figure 1.

In case of zirconium addition (Figure 1) the lattice parameter increases up to 1 wt % and then it tends to flatten up to 1.5 wt %. The solubility of zirconium is limited above 1 wt %. The peaks of monoclinic ZrO<sub>2</sub> were observed at 3 and 8 wt %. The amount of ZrO<sub>2</sub> entering the spinel lattice is not calculated at these weight percentages.

In case of tin the solid solubility is observed up to 5 wt %, the limit of this study. The lattice parameter increases monotonically for all weight percent addition of the dopant.

## 4.10 MAGNETIC MEASUREMENTS

## 4.10.1 Saturation magnetisation

The variation of magnetic moment with the applied field for base Ni 3Zn 7Fe2O4 at room temperature is given in Table 8 and illustrated in Figure 2. The maximum field applied is 10KOe. The variation of saturation magnetic moment at room temperature with the addition of the dopants is given in Table 9 and Figures 7,8 and 9.

In case of titanium addition the magnetic moment falls steeply at low percentages followed by an increase and again falls to low values at higher percentages as seen in Figure 7. The minimum moment of the initial dip is at 0.5 wt % of titanium addition.

With 0.25 wt % of zirconium addition the magnetic moment falls steeply from the initial value. Upto 1 wt% the moment decreases by a small amount and above 1 wt% it increases by a small amount (Figure 8). The initial steep decrease of moment is found in tin additions upto 0.5 wt%. Further additions upto 5 wt% reduces the moment continuously by a small amount as shown in Figure 9. 4.10.2 Curie temperature

The variation of magnetic moment at residual field (40 Oe) with temperature was studied for Ni Zn 7Fe<sub>2</sub>O<sub>4</sub> after 2,4 and 6 hours of firing duration. The samples were reground after each firing and the nature of the curve improved significantly with increased firing duration. The samples after 6 hours of firing shows a relatively sharper transition as seen in Figure 3.

In Figures 4,5 and 6 the temperature variation of the magnetic moment at residual field (40 0e) for titanium, zirconium and tin additions is shown. The curie temperatures (T<sub>c</sub>) of all samples were determined with the help of a computer program (provided by Dr. M. Misra) involving spline interpolation technique. The curie temperature of base Ni 3Zn 7Fe2O4 was compared with Misra and they match very well. Table 9 shows the curie temperature calculated for all samples doped with titanium, zirconium and tin.

Figure 7 shows the variation of curie temperature with the addition of titanium in Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$ . The transition temperature ( $T_c$ ) initially falls steeply and shows a minimum at 0.5 wt %. This is followed by an increase upto 1.5 wt % when again it starts lowering at 3 and 8 wt %. The trend of this curve is similar to that of the magnetic moment variation curve (Figure 7). Similar dip is found in zirconium additions with minimum  $T_c$  at 0.25 wt % (Figure 8). Two dips are found for tin additions at 0.25 and 1 wt % as shown in Figure 9. In both zirconium and tin additions the dips are not as prominent as in titanium.

#### CHAPTER 5

#### DISCUSSION

#### 5.1 SATURATION MAGNETISATION

The present observations of variation of magnetic moment when high valent cations such as titanium, zirconium and tin substituted Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$  are the following:

- There is a steep fall of initial magnetic moment in all three additions.
- 2) The initial steep fall is followed by a subsequent rise and fall in magnetic moment in the case of titanium addition (Fig. 7).
- 3) The initial steep fall is followed by a small rise in magnetic moment in the case of zirconium addition (Fig. 8).
- 4) The initial dip is followed by a gradual decrease of magnetic moment in the case of tin addition (Fig. 9).

The dip in magnetic moment when similar high valent cations were substituted in Ni-Zn ferrite was observed by earlier workers 8,26,27. The observed phenomenon of sharp decrease of magnetic moment is quite contrary to the expected monotonic decrease when high valent non-magnetic ions are substituted in Ni-Zn ferrite.

Possible explanations were given by earlier workers 26,27,8 to this observed dip in the magnetic moment. Das 26 presumed the variation of magnetic moment with the addition of all non-magnetic ions to behave in a similar manner as the addition of zinc around the composition point. The variation was arbitrarily taken from saturation magnetisation (M<sub>S</sub>) versus zinc content curve (Gorter 2) and calculated the variation of M<sub>S</sub> in titanium substituted Ni 5 To 5 Fe 204 based on Yafet-Kittel 28 relationship as,

$$M_s = M_b \sin \varphi - M_a$$

where  $M_b$  and  $M_a$  are the saturation magnetisation of B and A sublattices and  $(\pi-2\phi)$  is the canting angle of the sublattice components of B. The calculation of the slope  $\frac{dM}{dx}$ ,

where x is the amount of dopant added was shown to be negative both for A or B site substitution of titanium ion with slopes being different.

On the same basis Sen<sup>27</sup> calculated the slope of the magnetic moment with titanium addition in Ni<sub>3</sub>Zn<sub>.7</sub>Fe<sub>2</sub>O<sub>4</sub> and showed negative slopes in A or B substitutions. There was no explanation given by the above workers for the subsequent rise in the magnetic moment. But in both the cases a larger negative slope was seen in A site occupancy

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than for B site. In case of Das the slope was 1.6 times higher for A site occupancy and it was 5.7 times in the case of Sen. This difference may be due to two different compositions of the ferrite and hence two different canting angles.

The analysis of cation distribution and calculation of saturation magnetisation by Misra based on Satyamurthy's 3-sublattice model gives a better reasoning for the observed phenomenon. In the three sublattice model of ferrimagnetism the B sublattice is divided into  $B_1$  and  $B_2$  each of which is oriented at an angle  $(\pi-\phi)$  in opposite direction with respect to the moment of A sublattice where  $\phi$  is the Yafet-Kittel angle. He had calculated the canting angle for different additions of titanium from the interaction parameters of different cations which enabled him to calculate the cation distribution of titanium substituted Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$ .

Based on Misra's model, for small concentrations of titanium ions, the substituting ion enters the A site only.

After a critical value of concentration it starts entering B site and at higher concentrations it enter A and B sites with equal preference.

His model of  $Ti^{4+} \rightarrow Fe^{3+} + \frac{1}{2} (Ni_{3}Zn_{7})^{2+}$  resulting in half vacancy of divalent cation and  $Ti^{4+}$  going to A site together with some transfer of  $Fe^{3+}$  from A to B site accounts

for the initial dip. After a critical composition, Ti<sup>4+</sup> enters B site resulting in a transfer of Fe<sup>3+</sup> from B to A site, thereby increasing the magnetic moment. At higher percentages Ti<sup>4+</sup> enters both A and B sites and the transfer process of Fe<sup>3+</sup> is reduced to a low value. Similar analysis like Misra could not be carried out in the present work as the interaction parameters and the canting angle were not calculated. This restricted us in determining the cation distribution and hence the theoretical magnetic moment.

Misra's model may qualitatively explain the behaviour of  $M_S$  variation with dopant addition in the case of zirconium which shows a sharp initial fall followed by a small rise though this dip and subsequent rise is accentuated in the case of titanium. In the case of tin the initial steep fall is followed by a gradual slope of the curve.

This model is unable to explain the sharp initial dip in magnetic moment of 29.76 percent per ½ wt % of titanium addition in Misra's observation and 26.63 percent, 33.92 percent, 19 percent, per 1/4 wt % of titanium, zirconium and tin addition respectively in the present study. The weight percentage of titanium added was twice in the case of Misra with respect to the present observation for nearly the same initial fall of magnetic moment. Also, the initial dip and subsequent rise is more pronounced in the present work and it may be due to the better chemical homogenity of the dopant in the basic ferrite.

### 5.2 LATTICE PARAMETER

The following are the observation of lattice parameter variation with the addition of high valent cations such as titanium, zirconium and tin in Ni 3Zn 7Fe 204 in the present study.

- 1) The dip in lattice parameter was observed only in the case of titanium (Fig. 1). The dip is followed by a monotonic increase with higher additions of titanium.
- 2) The lattice parameter of zirconium and tin additions increases linearly as the ions presumably go into the solid solution (Fig. 1).
- The peaks of monoclinic ZrO<sub>2</sub> phase were observed at higher percentages indicating the limited solubility of zirconium.

The observations of lattice parameter variation with similar high valent cation addition in Ni-Zn ferrite by earlier workers 8,25 are the following:

- 1) A dip in lattice parameter was observed upto 0.8 wt titanium substituted Ni<sub>.5</sub>Zn<sub>.5</sub>Fe<sub>2</sub>O<sub>4</sub> followed by a monotonic rise by Das<sup>8</sup>.
- 2) A dip in lattice parameter was observed by Misra<sup>8</sup> for titanium substituted Ni<sub>3</sub>Zn<sub>7</sub>Fe<sub>2</sub>O<sub>4</sub> upto 1 wt % followed by a subsequent monotonic increase.

The linear increase of lattice parameter observed was followed by the flattening of the curve with zirconium, niobium additions in Ni.5<sup>Zn</sup>.5<sup>Fe</sup>2<sup>O</sup>4 showed the limit of their solubility in the case of Das<sup>8</sup>.

There was no appropriate explanation given for the initial dip observed for titanium additions by Das. Misra's explanation to the observed variation of lattice parameter of titanium substituted Ni<sub>3</sub>Zn<sub>.7</sub>Fe<sub>2</sub>O<sub>4</sub> was based on cation distribution resulting in an initial decrease of net charge at A site followed by a subsequent increase at later stage. From the analysis of cation distribution he concluded that

- 1) At very low percentages titanium enters the A site.
- 2) At higher percentages after a critical value, titanium starts entering the B site.
- 3) The entry of Ti4+ in the spinel lattice creates vacancies.

As a result, the charge at A site decreases with the entry of titanium, because of the combined effect of (a) excess charge of Ti<sup>4+</sup> over one Fe<sup>3+</sup> as it enters the site (b) transfer of same Fe<sup>3+</sup> from A to B site and (c) removal of half divalent ion (zinc valency) from the site to maintain the charge balance.

The lowering of charge at A site increases the Madelung constant (Gorter<sup>2</sup>) and the fact that the lattice parameter varies inversely with Madelung constant accounts for the initial dip. At higher percentages, Ti<sup>4+</sup> starts entering the

B sites which increases the charge at A site leading to the increase in lattice parameter. Also, with higher additions of titanium, vacancies increase which tend to repel the neighbouring ions thereby causing further increase in the lattice parameter.

The explanation based on Misra fails to account the absence of initial dip in other qquatrivalent cations such as zirconium and tin (Fig. 1) additions in Ni $_3$ Zn $_7$ Fe $_2$ O $_4$ . The substitutions by zirconium and tin increases the lattice parameter as they presumably enter the solid solutions. The present observations of lattice parameter variation with zirconium addition are similar to that of Das $^{25}$ .

In this case (Fig. 1) the limit of solubility of zirconium in Ni $_3$ Zn $_7$ Fe $_2$ O $_4$  is 1 wt %, after which the linear increase of lattice parameter tends to flatten out. Also, peaks of  ${\bf ZrO}_2$  were observed at 3 and 8 wt % indicating the limited solubility of zirconium. Das $^{25}$  had found the solubility limit to be 3.40 wt %  ${\bf ZrO}_2$  in Ni $_5$ Zn $_5$ Fe $_2$ O $_4$ .

# 5.3 CURIE TEMPERATURE

The nature of the curves of magnetic moment variation with temperature in Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$  improved with increase in firing time from 2 to 6 hours at  $_{1200}$ °C(Fig. 5). This clearly indicates the improvement in homogenisation of the dopant in the ferrite matrix and better ferritisation of the material.

The dips in curie temperature with the dopant addition are observed in the additions of titanium, zirconium and tin (Figs. 7,8 and 9). The dips are followed by a rise in all the dopant additions, the rise being maximum in the case of titanium.

Gorter<sup>2</sup> showed that the substitutions of now magnetic ions in tetrahedral site cause a decrease in curie temperature and the substitution in octahedral site also decreases the curie temperature but relatively less. The addition of zinc ferrite to nickel ferrite decreases the curie point. Expression for determining the curie temperature of Ni-Zn ferrite was given by Gilleo<sup>4</sup> involving the interaction parameters. But to a system like ours no theoretical calculation of curie temperature is available. Misra<sup>8</sup> had indicated the possibility of determining the curie temperature using a 5-sublattice model. The working out of the problem will be an extensive theoretical works.

The observed variations of curie temperature with the quatrivalent cations like titanium, zirconium and tin are similar (Figs. 7,8 and 9). The initial fall and rise of curie temperature is more prominent in the case titanium than in zirconium and tin additions. The observed variation is not understood in any of the above high-valent cation addition in Ni  $_3$ Zn  $_7$ Fe $_2$ O4.

The temperature variation of magnetic moment of the titanium, zirconium, and tin additions is shown in Figs. 4,5 and 6. In all the cases the magnetic moment does not become zero and a small tail is observed near the curie point. Our method of sample preparation is expected to be more chemically homogenous and this tail cannot be attributed to chemical inhomogenity alone. Further interpretation in terms of short-range ordering is required to explain the observed tail. However, no further analysis of the tail was carried out and curie temperature was determined from  $\frac{d\sigma}{dT}$  versus T curves.

#### CHAPTER 6

#### SUMMARY AND CONCLUSIONS

 $^{\rm Ni}$ . $^{\rm 7Fe}$ 2 $^{\rm O}$ 4 was doped with titanium, zirconium and tin by the solution method to a limit of 8 wt %, 8 wt % and 5 wt % respectively.

Step-scan method in conjunction with X-ray diffractometer was used for lattice parameter determination. The accurate extrapolated lattice parameter values were calculated from the lattice parameter versus Nelson-Riley function plot by least square fit of a straight line.

Vibrating magnetometer was used to determine the saturation magnetisation at 10KOe and the curie temperature at a residual field of 40 Oe.

For Ti<sup>4+</sup> doped samples, the lattice parameter drops sharply at low percentages followed by a monotonic rise upto 8 wt %. This may be explained on the basis of increase in charge due to the entry of Ti<sup>4+</sup> ions in the A sites at low percentages. The entry of Ti<sup>4+</sup> ions to B sites at higher percentages decreases the net-charge at A site and also the number of vacancies increase thereby further increasing the lattice parameter.

The lattice parameter of  $Zr^{4+}$  ions doped samples increases upto 1 wt % and then tends to flatten off indicating the limit of solubility of zirconium. The lattice parameter of  $Sn^{4+}$  increases monotonically upto the limit of the study. The lattice parameter increase in both  $Zr^{4+}$  and  $Sn^{4+}$  ions is explained on the basis of cation vacancies.

Sharp initial decrease of magnetic moment is observed with the additions of  ${\rm Ti}^{4+}$ ,  ${\rm Zr}^{4+}$ , and  ${\rm Sn}^{4+}$ . In the case of titanium addition the initial dip is followed by subsequent rise and again a fall in magnetic moment at higher percentages. In the case of zirconium addition the initial dip is followed by a gradual slope and a very small increase and with tin addition the initial dip is followed by a gradual change of slope.

The observed variations of magnetic moment in all the three cation additions may be explained based on the model of transfer process of Fe<sup>3+</sup> ions depending on the site preference of the substituting ion.

With the addition of titanium, zirconium and tin the curie temperature falls steeply at low percentages, followed by subsequent rise. This is maximum in the case of titanium. These observations could not be explained for any of the dopant addition although a similar variation has been observed earlier.

The refiring and regrinding of the samples improved the relative sharpness of the transition of magnetic moment with temperature. The samples fired for 6 hours with regrinding showed a relatively sharper transition indicating the better homogenity.

Table 1

Ionic radii and atomic weights of component elements\*

Elements	Atomic weight	Normal valency state	Ionic radius (OA)	
Nickel (Ni)	58.71	2 <sup>+</sup>	0.69	
Zinic (Zn)	65.37	2+	0.74	
Iron (Fe)	55,85	2 <sup>+</sup> 3 <sup>+</sup>	0.64	
Titanium (Ti)	47,90	4+	0.68	
Zirconium (Zr)	91.22	4 <sup>+</sup>	0.79	
Tin (Sn)	118.69	4 <sup>†</sup>	0.71	

<sup>\*</sup>Reference: - 'Hand Book of Chemistry and Physics', CRC Press, both edition, 1980-81, p. F-214.

Table 2
Batch composition

(i) Packing material

Component	Molecular weight	Composition (parts)	Percentage weight of the component	Grad <b>e</b>
NiO	74.71	15	9,375	L.R.
ZnO	81.37	35	23.825	L.R.
Fe <sub>2</sub> O <sub>3</sub>	159.70	50	66,800	L.R.

(ii) Base-ferrite (Ni\_3Zn\_7Fe2O4)

Metal	Atmonic weight	Percentage weight of the metal	Metal oxide formed	Percentage weight of metal oxide formed
lickel	58,71	10,06	NiO	9,37
Zinc	65.37	26.14	ZnO	23.83
[ron	55.85	63,80	Fe <sub>2</sub> O <sub>3</sub>	66,80

Dopants additions to the base solution

Table

Dopant	Mole fraction of the dopant	wt % of the dopant	t of opant	Volume of the dopant solution	Volume of the base solution	Strength of base-solution (cms/m1)
			(gms)	(m1)	(m1)	/ / <sub>2</sub> 6 )
Colors the state of the state o	2	(*)	4	2	9	7
	0,0075	0,25	0.0282	3,40	00.06	
	0,0150	0,50	0.0564	6,80	00.0	
Titanium	0.0300	1,00	0,1128	13,55	00.06	0.1255
	0,0450	13. 30.	0.1693	20.30	00°00	
	0060.0	3,00	0,3385	40.60	90,00	
	0.2390	8,00	0,9027	108,30	00°06	•

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	0.0049	0,25	0.0250	2,20	59.90		· 1110年 - 110年 - 1110年 - 1110
	0,0097	0.50	0.0500	4,40	59.75		
Zirconium	0.0194	1.00	0.1000	8,80	59,45	0.1655	
	0.0291	1,50	0.1500	13,20	59,15		
	0.0582	3,00	0.3000	26.25	58,25		
	0.1550	00°8	0,8000	70,00	55,25		
	0.0040	0,25	0,0250	1,25	29,90		
	0*0000	0,50	0,0500	2,50	59,75		
Tin	0,0160	1.00	0,1000	2,00	59,45	0.1655	
	0.0320	2,00	0,2000	10,00	58,85		
	0.0800	5,00	0.5000	25,00	57,05		

Table 4 X-ray diffraction lines Ni 3Zn 7Fe2O4 compared with ASTM data

ASTM X-ray hkl	data for Ni-Zn ferrito d A <sup>o</sup>	Observed dispacing for Ni.3 <sup>Zn</sup> .7 <sup>Fe</sup> 2 <sup>0</sup> 4
		B - California in the colorest of the second
111	4.85	eu
220	2.966	2.985
211	2.699	May
311	2.533	2,535
222	2.423	2.424
400	2.100	2.100
422	1.715	1.718
511,333	1.617	1.619
440	1,485	1.490
531	1.417	and .
620	1.327	1.328
583	1.280	1,283
444	1.212	Sed .
551,711	1.174	NA NA
642	1.122	_
553,731	1.093	. ·
800	1.043	*
X-ray unit Target Scan speed Chart speed	= Seifert (Germany) = Cromium = 30min. =150/min.	

contd...

Table 5

Estimation of lattics parameter of titanium doped  $^{\rm Ni}_{\rm *}{}_{\rm 3}^{\rm Zn}_{\rm *}7^{\rm Fe}{}_{\rm 2}^{\rm 0}{}_{\rm 4}$ 

	Extra- polated value of la- ttice para- meter	(A0)	8,4178
	Slope	6	0.0061
	Nelson- Riley fun- ctions f( $\Theta$ )	æ	2.8804 1.9041 1.7565 1.1917 0.6325 0.5056 0.3502 0.1247
	Lattice para- meter for di- fferent d values (Ao)		8.4010 8.4038 8.4041 8.4016 8.4210 8.4210 8.4294 8.4294 8.4155
	hkl	9	220 311 222 400 422 511 440 620 533
	d (Ao)	ಬ	2.9702 2.5354 2.4260 2.1004 1.7189 1.6200 1.4901 1.3286
The state of the s	G de- Crees)	4	22,685 26,860 28,175 33,050 41,790 45,000 50,240 50,240
	2 <del>0</del> (de- grees)	3	45,370 3,720 56,350 66,100 83,580 90,000 100,480 119,130
The second secon	Mole fract- ion	2	
And the state of t	Weight percen- tage	F.	

															4	14		
10			8,4127				-						8,4160		Ĭ.			
6			-0.0056				. *						-0,0019					contd
æ	2,8797	1,7529	1,1935	0,6291	0,5038	0,3468	0,1753	0,1295	2,8882	1,9650	1,7650	1.1940	0,6302	0,5050	0,3471	0,1773	0,1295	
7	8,4001	8,3973	8,4061	8,4083	8,4111	8,4129	8,4061	8,4174	8,4107	8,4103	8,4199	8,4073	8,4124	8,4155	8,4145	8,4184	8,4174	
9	220	222	400	422	511	440	620	533	220	311	222	400	422	511	440	620	533	
5	2,9699	2,4241	2,1015	1,7163	1,6187	1,4872	1,3291	1,2836	2,9736	2,5358	2,4306	2,1018	1,7171	1,6197	1,4875	1,3311	1,2836	
4	22_687	28,20C	33,030	41,867	45,045	50,375	59,525	63,175	22,657	26,855	28,117	33,025	41,842	45,015	50,362	59,382	63,175	
င	45,375	56,400	090*99	83, 735	060.06	100,750	119,050	126,350	45,315	53,710	56,235	090*99	83,685	90,030	100,725	118,765	126,350	
2			0 0075									0.0150						
터			0.05	•								0,50						

10			8,4153							8,4162					
6			-0.0046							-0.0087					
ω	2,8832	1.7602	0.6288	0.3463	0,1760	0.1310	2,8720	1,9562	1,7580	1,1940	0,6290	0,5038	0.3463	0,1770	0.1295
7	8,4045 8,4008	8,4110	8,4071	8,4103 8,4105	8,4104	8,4278	8,3905	8,3951	8,4068	8,4073	8,4079	8,4103	8,4105	8,4163	8,4174
9	220 311	222	422	511	620	533	220	311	222	400	422	511	440	620	533
ව	2,9714	2,4280	1.7161	1.0186 1.4869	1,3298	1,2852	2,9665	2,5312	2,4268	2,1018	1,7163	1,6186	1,4868	1,3307	1,2836
4	22,675 26,887	28,150 33,010	41,875	45,050 50,395	59,475	63,035	22,715	26,907	28,165	33,025	41,870	45,050	50,395	59,407	63,175
3	45,350	56,300	83,750	100, 790	118,950	126,070	45,430	53,815	56,330	090*99	83,740	90,100	100,790	118,815	126,350
2		000 0 0300			•					0.0450					
Ħ		1,00								1.50	•				

contd...

T. Marriage Confess and Transfer Confession					δ			,						Ω				
Particular of the second	10				8,4178									8,4205				
	6				-0.0049									-0,0023				
The state of the s	ω .	2,8832	1,9616	1,1978	0,6306	0,5047	0,3475	0.1771	0,1295	2,8903	1,9684	1,7639	1,1985	0,6321	0,5054	0,3481	0,1780	0,1298
	7	8,4045	8,4045	8,4169	8,4136	8,4140	8,4160	8,4170	8,4174	8,4133	8,4161	8,4178	8,4186	8,4194	8,4170	8,4190	8,4228	8,4192
Service of the Service of	, 9	220	311	400	422	511	440	620	533	220	311	222	400	422	511	440	620	533
	വ	2,9714	2,5340	2,1042	1,7174	1,6193	1,4878	1,3308	1,2836	2,9746	2,5375	2,4300	2,1046	1,7186	1,6199	1,4883	1,3318	1,2839
	4	22,675	26,875	32,982	41,835	45,025	50,350	59,400	63,175	22,650	26,835	28,125	32,975	41,800	45,005	50,325	59,332	63,150
	က	45,350	53,750	65,965	83,670	90,050	100,700	118,800	126,350	45,300	53,670	56,250	65,950	83,600	90.010	100,650	118,665	126,300
	2	*			0060 0 00									0,2390				
					0		-							8,00			•	

Estimation of lattice parameter of zirconium doped  $^{\rm Ni}_{\rm \bullet} \, ^{\rm Z^{\rm o}}_{\rm \bullet} \, ^{\rm 7^{\rm Fe}}_{\rm 2^{\rm o}}_{\rm 4}$ Table 6

Weight percen- tage	Mole fra- - ction	20 (de- grees)	de- grees)	(AQ)	hk1	Lattice parameter for di- fferent d values (Ao)	Nelson r Riley fun- ction f(0)	Slope	Extrapo- lated value of lattice parameter (Ao)
Н	8	ε	4	2	9	7	ω	6	10
		45,380	22,690	2,9696	220	8 3992	2,8790		
		53,740	26,870	2,5345	311	8,4057	1,9625		
		26,330	28,165	2,4268	222	8,4068	1,7580		
		66,040	33,020	2,1021	400	8,4084	1,1944		Ì
0.25	0.0049	83,660	41,830	1,7176	422	8,4145	0,6308	9900.0-	8,4183
		90.035	45,017	1,6195	511	8,4151	0,5049		
		100,700	50,350	1,4877	440	8,4160	0,3475		
		118,800	59,400	1,3308	620	8,4169	0,1771		
		126,335	63,167	1,2837	533	8,4179	0.1296	-3	3 3

contd.

1																		
10					8,4196									8,4214			,	
6					-0.0033									-0.0027				
8	2,8882	1,9672	1,7610	1,1967	0,6314	0,5059	0,3480.	0,1774	0,1298	2,8903	1,9680	1,7639	1,1983	0,6318	0,5070	0,3482	0.1778	0,1298
7	8,4107	8,4139	8,4123	8,4141	8,4169	8,4192	8,4190	8,4189	8,4192	8,4133	8,4153	8,4178	8,4180	8,4181	8,4236	8,4196	8,4217	8,4192
9	220	311	222	400	422	511	440	620	533	220	311	222	8	422	511	440	620	533
5.	2,9736	2,5369	2,4284	2,1035	1,7181	1,6203	1,4883	1,3311	1,2839	2,9746	2,5373	2,4300	2,1045	1,7183	1,6211	1,4884	1,3316	1,2839
4	22,657	26,842	28,145	32,995	41,815	44.990	50,325	59,377	63,150	22,650	26,837	28,125	32,977	41,807	44,960	50,320	59,345	63,150
က	45,315	53,685	56,290	65,990	83,630	89,980	100,650	118,755	126,300	45,300	53,675	56,250	65,955	83,615	89,920	100,640	118,690	126,300
2					7600.0							•		0.0194				
F			1 *		0.50									1.00		*		

	1							•								•		
														•				
10					8,4203	•								4100				
6					-0.0024			* .						9650		•		
8	2,8903	1,9684	1,7628	1,1981	0,6317	0,5060	0,3483	0,1776	0.1298	0300	1 0557	1.7551	1.1924	0.6284	0,5030	0.3454	0.1753	0.1300
7	8,4133	8,4160	8,4158	8,4175	8,4177	8,4195	8,4199	8,4200	8,4198	2000	8 3943	8,4014	8,4033	8,4054	8,4071	8,4060	8,4061	8,4211
9	220	311	222	400	422	511	440	620	533	000	311	222	90	422	511	440	620	533
5.	2,9746	2,5375	2,4294	2,1044	1,7183	1,6203	1,4884	1,3313	1.2840	2,9696	2,5310	2,4253	2,1008	1,7158	1,6179	1,4860	1,3291	1,2842
4	22,650	26,835	28,132	32,980	41,810	44,987	50,317	59,365	63,142	22,690	26,910	28,185	33,042	41,835	45,072	50,432	59,525	63,125
ო	45,300	53,670	56,265	096*59	83,620	89,975	100,635	118,730	126,285	45,380	53,820	56,370	66,085	83,770	90,145	100,865	119,050	126.760
7					0.0291									0.0582			•	
					යි .									8	•			

9 10				-0.0023 8.4179					
							•		*
ω	2,8889	1,9659	1,7613	1,1972	0,6310	0,5063	0.3477	0.1772	0.1292
7	8,4115	8,4117	8,4130	8,4152	8,4153	8,4206	8,4172	8,4176	8,4155
9	220	311	222	400	422	511	440	620	533
2	2,9739	2,5362	2,4286	2,1038	1,7178	1,6206	1,4880	1,3309	1,2833
4	22,655	26,850	28,142	32,990	41,825	44,980	50,340	59,392	63,200
ဇ	45,310	53,700	56,285	65,980	83,650	096.68	100,680	118,785	126,400
2				8,00 0,1550					
T				8,00			1 m v		

contd ...

Table 8 Estimation of lattice parameter of tin doped  $${\rm Ni}_{\mbox{\scriptsize o}}{\rm 3}^{\rm Z}{\rm n}_{\mbox{\scriptsize o}}{\rm 7}^{\rm Fe}{\rm 2}^{\rm O}{\rm 4}$ 

								The second second second second	
82	% Mole- fraction	20 n (de- grees)	0 (de-	(Ag	hk1	Lattice para- meters for different d-values(A)	Nelson- Riley function f(0)	Slope	Extrapola- ted value of lattice para- meter(A)
	2	က	4	ک	9	7	ω	6	10
		45,275	23,637	2,9761	220	8,4177	2,8938		
		53,675	26,837	2,5373	311	8,4153	1,9680		
		56,250	28,125	2,4300	222	8,4178	1,7639		
		65,950	32,975	2,1046	9	8,4185	1,1985		
.25	0.0040	83,590	41,795	1,7188	422	8,4202	0,6323	0,0020	8,4215
		89,945	44,972	1,6208	511	8,4217	0,5065		
		100,615	50,307	1,4887	440	8,4212	0,3485		
	L/n	118,725	59,362	1,3313	620	8,4202	0,1776		
	1 /1	126,215	63,107	1,2844	533	8,4223	0.1302		
	À								

			8,4215	•									8,4287					•
			-0,0062 8							*			<b>-0.0042</b> 8.	in the second				contd.
2,9846	1,9646	1,7620	1,1964	0,6322	0,5063	0,3482	0.1781	0,1303	2,8924	1,9718	1,7661	1,1985	0.6359	0,5079	0.3496	0,1787	0.1307	
8,4063	8,4095	8,4144	8,4135	8,4198	8,4206	8,4196	8,4234	8,4229	8,4159	8,4218	8,4219	8,4186	8,4334	8,4273	8,4267	8,4269	8,4256	
220	311	222	400	422	511	440	620	533	220	311	222	400	422	511	440	620	533	
2,9720	2,5356	2,4290	2,1033	1,7187	1,6207	1,4884	1,3319	1,2845	2,9755	2,5393	2,4312	2,4046	1,7214	1,6218	1,4896	1,3324	1,2848	
22,670	26,857	28,137	32,997	41,797	44,980	50,320	59,325	63,100	22,642	21,815	28,110	32,975	41,715	44,935	50,262	59,285	63,065	
45,340	53,715	56,275	65,995	83,595	89,960	100,640	118,650	126,200	45,285	53,630	56,220	65,950	83,430	89.870	100,525	118,570	126,130	
			0.0080										0.0160					
			200									•	1.00	•				

 $\infty$ 

က

N

10					8,4288				•				•	8,4376					
6					-0.0074									-0.0055					
ω	2,8368	1,9676	1,7624	1,1981	0.6330	0.5076	0,3499	0.1784	0,1314	2,9024	1,9659	1,7706	1,2045	0,6365	0,5100	0,3524	0,1803	0,1317.	The second secon
7	8,4089	8,4146	8,4150	8,4174	8,4227	8,4261	8,4282	8,4252	8,4304	8,4283	8,4117	8,4302	8,4333	8,4358	8,4361	8,4404	8,4366	8,4323	
9	220	311	222	400	422	511	440	620	533	220	311	222	400	422	511	440	622	533	
S	2,9729	2,5371	2,4292	2,1043	1,7192	1,6216	1,4899	1,3321	1,2856	2,9798	2,5362	2,4336	2,1083	1,7219	1,6235	1,4920	1,3339	1,2859	
4	23,662	26,840	28,135	32,980	41,780	44,942	50,250	59,305	63,000	22,607	26,850	28,080	32,910	41,700	44,875	50,150	59,175	62,975	
8	45,325	53,280	56,270	65,960	83,560	89,885	100,500	118,610	126,000	45,215	53,700	56,160	65,820	83,400	89,750	100,300	118,350	125,950	
7					0.0320		٠		* * *					0080.0					
1					2.00									2.00					

Variation of magnetic moment of Ni 3<sup>Zn</sup> 7<sup>Fe</sup>2<sup>O</sup>4 with applied field at room temperature

Applied field (K. Oe)	Magnetic	moment
0.00	0.0007	0.098
0.1	0.0546	7.668
0.3	0,1570	22.050
0.5	0,2360	33.146
0.7	0,2920	41.011
1.0	0,3440	48.314
2.0	0,3970	55 <b>.75</b> 8
3.0	0.4110	57.724
4.0	0,4190	58,848
5.0	0.4240	59,550
6.0	0,4280	60.112
7.0	0,4310	60,533
8,0	0,4340	60,955
9.0	0,4370	61.376
10.0	0.4390	61.657

Room temperature = 24°C

wass of the sample = 0.0712 gms.

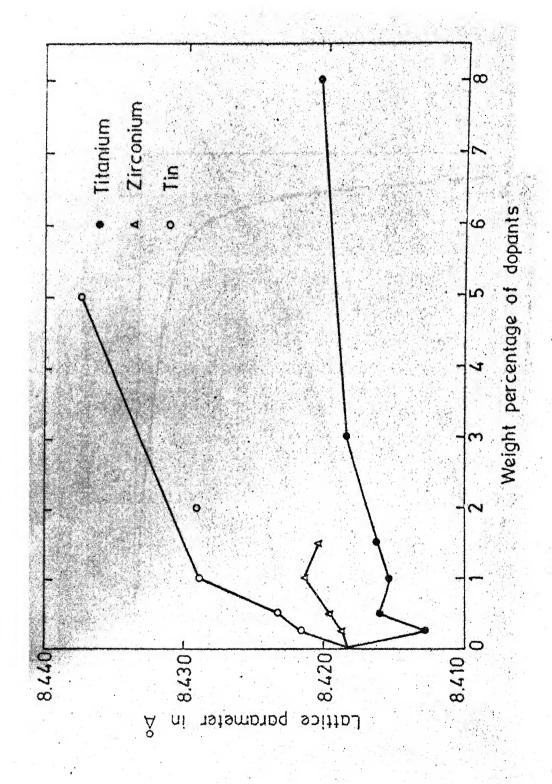
contd

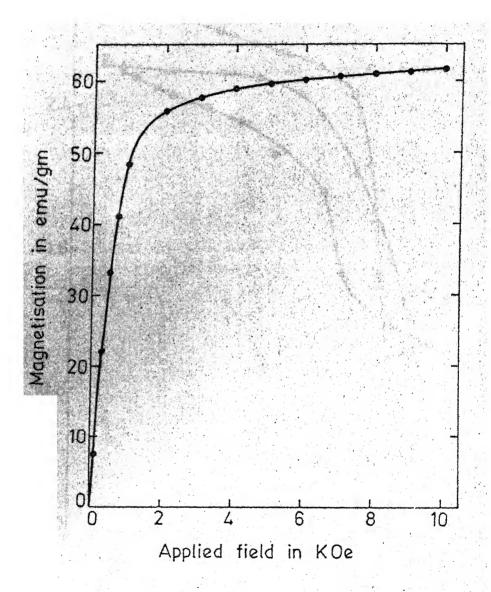
Table 9

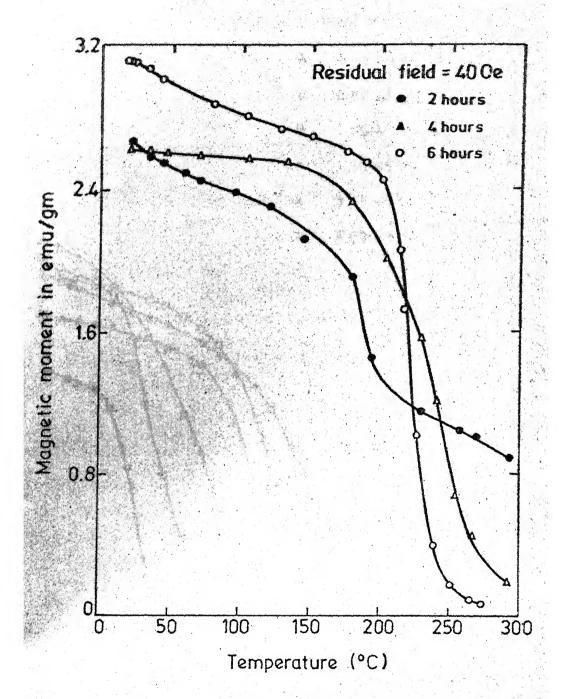
Saturations magnetic moment, curie temperature and lattice parameter of Ni  $_3$ Zn  $_7^{\rm Fe}{}_2{}^0{}_4$  with different dopants

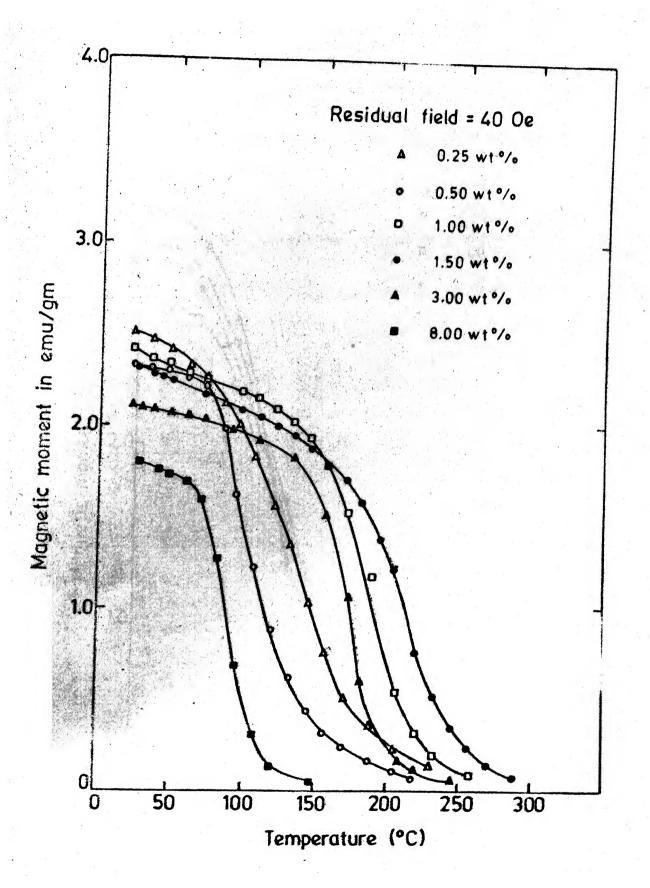
	wt% of the dopant	Mors of the Copant	Saturation magnetisation (emu/gm)	4π <sup>M</sup> (Gaßss)	Curie tempe- rature (°C)	Lattice parameter
1	2	ന	4	S	9	7
			61,657	4125.1	204 1	0 44.70
	0,25	0,0075	45,317	3031,9	129.8	0.41/8 8 7107
	0.50	0,0150	40,873	2734.5	112.8	0 + 1 Z < 8 × 1 / C O
	1,00	0,0300	44.046	2946.8	167.2	8 7752
Titanium	1,50	0,0450	53,953	3609.6	188.2	8 7160
	3.00	0060*0	48,168	3222,6	166.0	0 + 10v
	8,00	0,2390	28,961	1937,6	123.2	8,4205

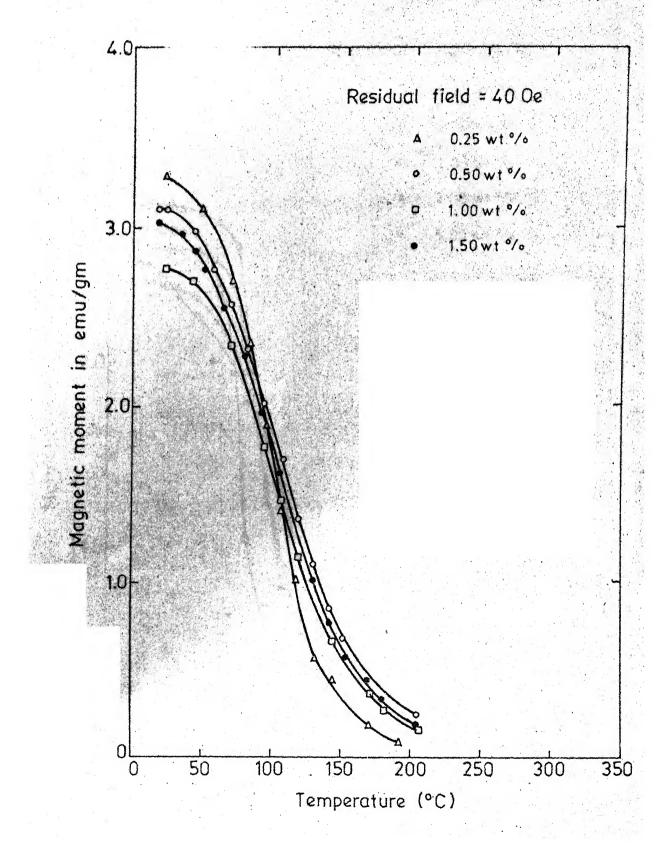
Andread Andread Control of the Contr	2	က	4	ប	9	7
*	0.25	0,0049	40,746	2726.0	101.4	8,4183
	0,50	0,0097	39,367	X 33,8	112,4	8,4196
	1,00	0.0194	38*666	2586,9	110.0	8,4203
Zirconium	1,50	0,0291	39,575	2647,7	109,8	8,4203
	3°00	0,0582	44,854	3000.9	216,8	8,4108
	8,00	0,1550	42,325	2831,7	144.9	8,4179
	0.25	0,0040	49,944	3341,4	107.0	8,4215
	0.50	0,0080	45,911	3071.6	113,8	8, 4232
	1,00	0,0160	44,690	2989.9	95.9	8,4287
Tin	2.00	0,0320	43,478	2908,8	110.9	8,4288
	2,00	0080	43,120	2884.9	113,9	8,4375

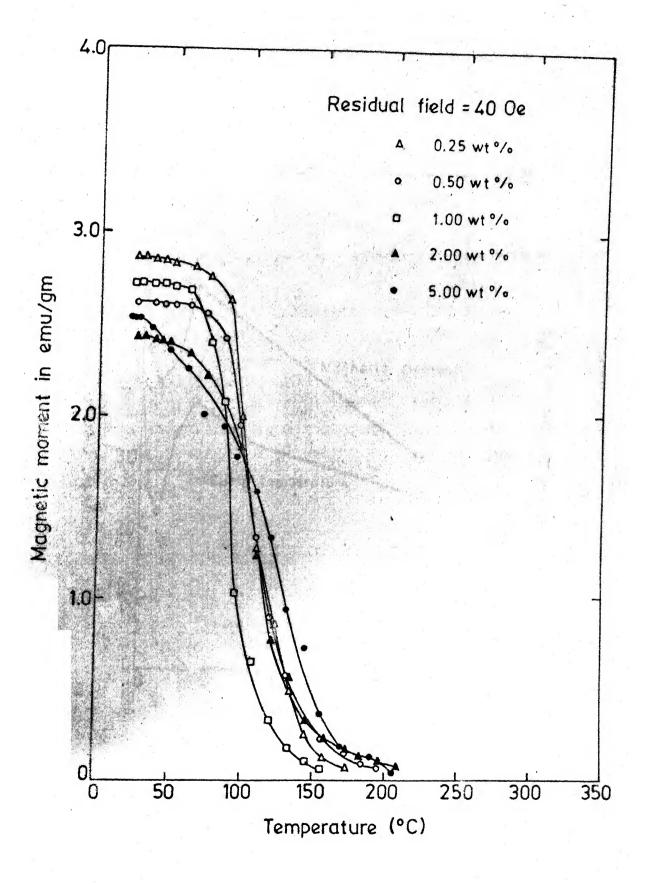


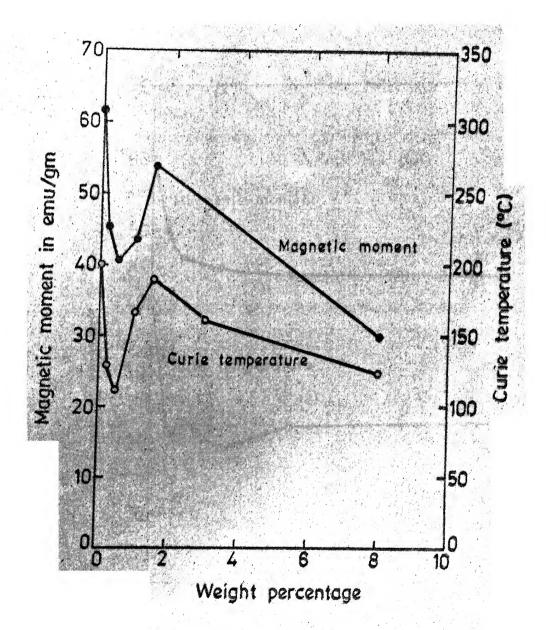


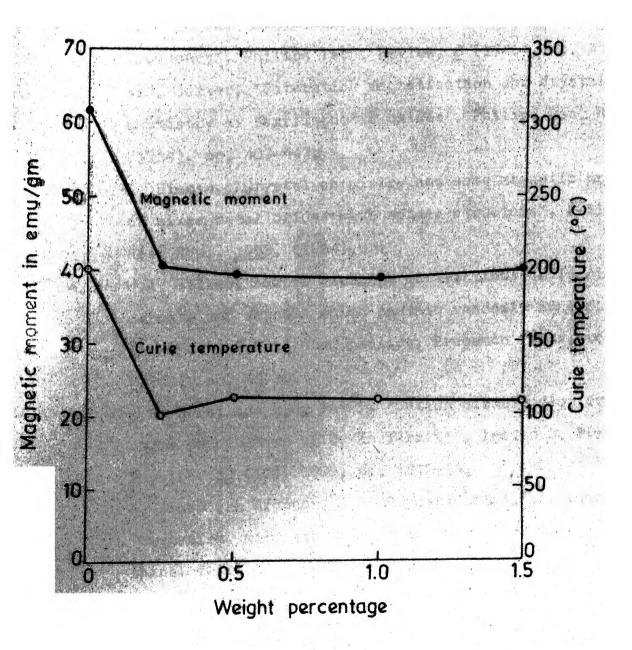












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### APPENDIX

### ESTIMATION OF LATTICE PARAMETER

The lattice parameter was precisely estimated using Nelson-Riley function.

The d values were calculated using Bragg's law

$$\lambda = 2d \sin\theta$$

where '20' was the diffraction angle and ' 'was the wave-length of chromium  $K_{\alpha}$  target.

The d-values calculated were used to find out the lattice parameter, 'a' as

$$a = d(h^2 + k^2 + 1^2)^{\frac{1}{2}}$$

where hkl were the indices for the diffracting planes adopted from ASTM data file for Ni-Zn ferrite.

Precise lattice parameter was obtained by extrapolating the lattice parameter to  $\Theta=90^{\circ}$  and using Nelson-Riley extrapolation function  $f(\Theta)$ . This function was chosen to use the low angle reflections for extrapolation of the lattice parameters. The Nelson-Riley function is given by

$$f(\theta) = (\frac{\cos^2\theta}{\sin\theta} + \frac{\cos^2\theta}{\theta})$$

Accurate determination of angle was done by step scan method. The X-ray counts were plotted against 20. The curves for 1 wt  $\cdot$  titanium substituted Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$  are shown in Fig. 10. The exact 20 corresponds to the maximum peak position which was measured accurately upto  $\pm$  0.01 degree.

A computer program was written to calculate the extrapolated lattice parameter from the lattice parameter versus Nelson-Riley function plot by least square fit of a straight line. The calculated d, hkl, lattice parameters and the corresponding Nelson-Riley function are given in Tables 5,6 and 7. The extrapolated lattice parameters are given in Tables 5,6,7 and 9, for titanium, zirconium and tin additions in Ni  $_3$ Zn  $_7$ Fe $_2$ O $_4$ .

